

(b) Energy to move a charge q from x_1 to x_2 is change in potential energy $\Delta(q\phi) = q(\phi_2 - \phi_1)$. Let $\phi_1 = 0$ at $x = \pm d$ and $\phi_2 = (1/2\epsilon_0)nqd^2$ at $x = 0$. Then

$$\mathcal{E} = \frac{1}{2\epsilon_0}nq^2d^2$$

Let $d = \lambda_D$; then

$$\mathcal{E} = \frac{1}{2\epsilon_0}nq^2 \frac{KT\epsilon_0}{nq^2} = \frac{1}{2}KT = E_{AV}$$

for a one-dimensional Maxwellian distribution. Hence, if $d > \lambda_D$, $\mathcal{E} > E_{AV}$. If the velocities are distributed in three dimensions, we have $E_{AV} = \frac{3}{2}KT$ and $\mathcal{E} > \frac{2}{3}E_{AV}$. The factor 3 is not important here. The point is that a thermal particle would not have enough energy to go very far in a plasma ($d > \lambda_D$) if the charge of one species is not neutralized by another species.

1-7. (a) $\lambda_D = 7400(2/10^{16})^{1/2} = 10^{-4}$ m, $N_D = 4.8 \times 10^4$.

(b) $\lambda_D = 7400(0.1/10^{17})^{1/2} = 2.3 \times 10^{-3}$ m, $N_D = 5.4 \times 10^4$.

(c) $\lambda_D = 7400(800/10^{17})^{1/2} = 6.6 \times 10^{-3}$ m, $N_D = 1.2 \times 10^5$.

2-1. $E = \frac{1}{2}mv_{\perp}^2 \therefore v_{\perp} = (2E/m)^{1/2}$, $r_L = mv_{\perp}/eB$.

(a)

$$v_{\perp} = \left[\frac{(2)(10^4)(1.6 \times 10^{-19})}{9.11 \times 10^{-31}} \right]^{1/2} = 5.93 \times 10^7 \text{ m/sec}$$

$$r_L = \frac{(9.11 \times 10^{-31})(5.93 \times 10^7)}{(1.6 \times 10^{-19})(0.5 \times 10^{-4})} = 6.75 \text{ m}$$

(b)

$$v_{\perp} = (300)(1000) = 3 \times 10^5 \text{ m/sec}$$

$$r_L = \frac{(1.67 \times 10^{-27})(3 \times 10^5)}{(1.6 \times 10^{-19})(5 \times 10^{-5})} = 6.26 \times 10^3 \text{ m} = 6.26 \text{ km}$$

(c)

$$v_{\perp} = \left[\frac{(2)(10^3)(1.6 \times 10^{-19})}{(4)(1.67 \times 10^{-27})} \right]^{1/2} = 2.19 \times 10^5 \text{ m/sec}$$

$$r_L = \frac{(4)(1.67 \times 10^{-27})(2.19 \times 10^5)}{(1.6 \times 10^{-19})(5.00 \times 10^{-5})} = 0.183 \text{ m}$$

(d)

$$r_L = \frac{2ME}{qB} = \frac{[(2)(4)(1.67 \times 10^{-27})(3.5 \times 10^5)(1.6 \times 10^{-19})]^{1/2}}{(2)(1.6 \times 10^{-19})(8)} \\ = 5.58 \times 10^{-3} \text{ m}$$

Plasma Physics And Controlled Fusion Solution Manual

DJ Losen



Plasma Physics And Controlled Fusion Solution Manual:

Solutions Manual for Controlled Fusion and Plasma Physics Miyamoto Kenro, Miyamoto Kenro Staff, 2006-07

Introduction to Plasma Physics and Controlled Fusion Francis F. Chen, 2013-03-09 TO THE SECOND EDITION In the nine years since this book was first written rapid progress has been made scientifically in nuclear fusion space physics and nonlinear plasma theory At the same time the energy shortage on the one hand and the exploration of Jupiter and Saturn on the other have increased the national awareness of the important applications of plasma physics to energy production and to the understanding of our space environment In magnetic confinement fusion this period has seen the attainment 13 of a Lawson number nTE of $2 \times 10^{21} \text{ cm}^{-3} \text{ sec}$ in the Alcator tokamaks at MIT neutral beam heating of the PL T tokamak at Princeton to KT_i 6.5 keV increase of average to 3% 5% in tokamaks at Oak Ridge and General Atomic and the stabilization of mirror confined plasmas at Livermore together with injection of ion current to near field reversal conditions in the 2XII device Invention of the tandem mirror has given magnetic confinement a new and exciting dimension New ideas have emerged such as the compact torus surface field devices and the E T mirror torus hybrid and some old ideas such as the stellarator and the reversed field pinch have been revived Radiofrequency heating has become a new star with its promise of dc current drive Perhaps most importantly great progress has been made in the understanding of the MHD behavior of toroidal plasmas tearing modes magnetic VII VIII islands and disruptions

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