

#1. $H_0 = -\frac{\hbar^2}{2m} \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right) + \frac{1}{2} m \omega^2 (x^2 + y^2)$

a) Is like a pair of harmonic oscillators, one in x , one in y .
For a single oscillator, lowest two states are [Eq. 2.49]:

$$\psi_0(x) = \left(\frac{m\omega}{\pi\hbar} \right)^{1/4} e^{-\frac{m\omega}{2\hbar} x^2} \quad E = \frac{1}{2} \hbar\omega$$

$$\psi_1(x) = \left(\frac{m\omega}{\pi\hbar} \right)^{1/4} \sqrt{2} \sqrt{\frac{m\omega}{\hbar}} x e^{-\frac{m\omega}{2\hbar} x^2} \quad E = \frac{3}{2} \hbar\omega \quad \text{similar for } \psi_n(y).$$

Lowest state of 2d oscillator is then just a product

$$\Psi_0(x, y) = \psi_0(x) \psi_0(y) = \left(\frac{m\omega}{\pi\hbar} \right)^{1/2} e^{-\frac{m\omega}{2\hbar} (x^2 + y^2)} \quad E_0 = \frac{\hbar\omega}{2} + \frac{\hbar\omega}{2} = \hbar\omega$$

Lowest excited states: put one quantum, either $n_x = 1$ or $n_y = 1$.

$$\Psi_1(x, y) = \psi_1(x) \psi_0(y) = \left(\frac{m\omega}{\pi\hbar} \right)^{1/2} \sqrt{2} \sqrt{\frac{m\omega}{\hbar}} x e^{-\frac{m\omega}{2\hbar} (x^2 + y^2)} \quad E_1 = \frac{\hbar\omega}{2} + \frac{3\hbar\omega}{2} = 2\hbar\omega$$

$$\Psi_2(x, y) = \psi_0(x) \psi_1(y) = \left(\frac{m\omega}{\pi\hbar} \right)^{1/2} \sqrt{2} \sqrt{\frac{m\omega}{\hbar}} y e^{-\frac{m\omega}{2\hbar} (x^2 + y^2)} \quad E_2 = 2\hbar\omega$$

b) $H' = \frac{b}{2} m \omega^2 xy (x^2 + y^2)$

Ground state is non-degenerate, apply usual 1st order formula:

$$E_0' = \langle \Psi_0 | H' | \Psi_0 \rangle = \frac{b}{2} m \omega^2 \left(\frac{m\omega}{\pi\hbar} \right) \int_{-\infty}^{\infty} dx \int_{-\infty}^{\infty} dy xy (x^2 + y^2) e^{-\frac{2m\omega}{\hbar} (x^2 + y^2)} = 0$$

Zero because $x > 0$ and $x < 0$ contributions are equal and opposite.

States 1 and 2 require degenerate perturbation theory.

Matrix of H' in this subspace involve matrix elements:

$$W_{ij} = \langle \Psi_i | H' | \Psi_j \rangle \quad \text{for } i, j = 1, 2 \quad W_{11} = W_{22} = 0 \quad \text{by inspection.}$$

$$\begin{aligned} W_{12} = W_{21} &= \frac{b}{2} m \omega^2 \frac{2}{\pi} \left(\frac{m\omega}{\hbar} \right)^2 \int_{-\infty}^{\infty} dx \int_{-\infty}^{\infty} dy xy (x^2 + y^2) xy e^{-\frac{2m\omega}{\hbar} (x^2 + y^2)} \\ &= \frac{b \hbar \omega}{\pi} \left(\frac{m\omega}{\hbar} \right)^3 \underbrace{\int_0^{2\pi} d\phi \left(\frac{1}{2} \sin 2\phi \right)^2}_{\frac{1}{4} \cdot \pi} \underbrace{\int_0^{\infty} dr \cdot r \cdot r^6 e^{-\alpha r^2}}_{\frac{d}{d\alpha} \int_0^{\infty} dr \cdot r e^{-\alpha r^2} = -\frac{3}{\alpha^4} = 3 \left(\frac{\hbar}{2m\omega} \right)^4}, \quad \alpha \equiv \frac{2m\omega}{\hbar} \end{aligned}$$

$$W_{12} = W_{21} = \frac{b \hbar \omega}{\pi} \left(\frac{m\omega}{\hbar} \right)^3 \cdot \frac{\pi}{4} \cdot 3 \cdot \left(\frac{\hbar}{2m\omega} \right)^4 = \frac{3}{64} b \cdot \hbar \omega \cdot \left(\frac{\hbar}{m\omega} \right).$$

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