Chapter One

1. $AC\{D,B\} = ACDB + ACBD$, $A\{C,B\}D = ACBD + ABCD$, $C\{D,A\}B = CDAB +$ CADB, and $\{C, A\}DB = CADB + ACDB$. Therefore $-AC\{D, B\} + A\{C, B\}D - C\{D, A\}B +$ $\{C, A\}DB = -ACDB + ABCD - CDAB + ACDB = ABCD - CDAB = [AB, CD]$

In preparing this solution manual, I have realized that problems 2 and 3 in are misplaced in this chapter. They belong in Chapter Three. The Pauli matrices are not even defined in Chapter One, nor is the math used in previous solution manual. – Jim Napolitano

2. (a) $Tr(X) = a_0 Tr(1) + \sum_{\ell} Tr(\sigma_{\ell})a_{\ell} = 2a_0 \text{ since } Tr(\sigma_{\ell}) = 0$. Also $\begin{array}{l} \operatorname{Tr}(\sigma_k X) = a_0 \operatorname{Tr}(\sigma_k) + \sum_{\ell} \operatorname{Tr}(\sigma_k \sigma_\ell) a_\ell = \frac{1}{2} \sum_{\ell} \operatorname{Tr}(\sigma_k \sigma_\ell + \sigma_\ell \sigma_k) a_\ell = \sum_{\ell} \delta_{k\ell} \operatorname{Tr}(1) a_\ell = 2 a_k. \text{ So, } \\ a_0 = \frac{1}{2} \operatorname{Tr}(X) \text{ and } a_k = \frac{1}{2} \operatorname{Tr}(\sigma_k X). \text{ (b) Just do the algebra to find } a_0 = (X_{11} + X_{22})/2, \\ a_1 = (X_{12} + X_{21})/2, \ a_2 = i(-X_{21} + X_{12})/2, \ \text{and } a_3 = (X_{11} - X_{22})/2. \end{array}$

 Since det(σ · a) = −a_x² − (a_x² + a_y²) = −|a|², the cognoscenti realize that this problem really has to do with rotation operators. From this result, and (3.2.44), we write

$$\det \left[\exp \left(\pm \frac{i \boldsymbol{\sigma} \cdot \hat{\mathbf{n}} \boldsymbol{\phi}}{2} \right) \right] = \cos \left(\frac{\boldsymbol{\phi}}{2} \right) \pm i \sin \left(\frac{\boldsymbol{\phi}}{2} \right)$$

and multiplying out determinants makes it clear that $det(\sigma \cdot \mathbf{a}') = det(\sigma \cdot \mathbf{a})$. Similarly, use (3.2.44) to explicitly write out the matrix σ - a' and equate the elements to those of σ - a. With $\hat{\mathbf{n}}$ in the z-direction, it is clear that we have just performed a rotation (of the spin vector) through the angle ϕ .

 (a) Tr(XY) ≡ ∑_a⟨a|XY|a⟩ = ∑_a∑_b⟨a|X|b⟩⟨b|Y|a⟩ by inserting the identity operator. Then commute and reverse, so Tr(XY) = ∑_b∑_a⟨b|Y|a⟩⟨a|X|b⟩ = ∑_b⟨b|YX|b⟩ = Tr(YX).
 (b) XY|α⟩ = X[Y|α⟩] is dual to ⟨α|(XY)†, but Y|α⟩ ≡ |β⟩ is dual to ⟨α|Y† ≡ ⟨β| and X|β⟩ is dual to $(\beta | X^{\dagger})$ so that $X[Y|\alpha)$ is dual to $(\alpha | Y^{\dagger}X^{\dagger})$. Therefore $(XY)^{\dagger} = Y^{\dagger}X^{\dagger}$ (c) $\exp[if(A)] = \sum_a \exp[if(A)]|a\rangle\langle a| = \sum_a \exp[if(a)]|a\rangle\langle a|$ (d) $\sum_a \psi_a^*(\mathbf{x}')\psi_a(\mathbf{x}'') = \sum_a \langle \mathbf{x}'|a\rangle^*(\mathbf{x}''|a) = \sum_a \langle \mathbf{x}''|a\rangle\langle a|\mathbf{x}'\rangle = \langle \mathbf{x}''|\mathbf{x}'\rangle = \delta(\mathbf{x}'' - \mathbf{x}')$

 For basis kets |a_i⟩, matrix elements of X ≡ |α⟩(β| are X_{ij} = ⟨a_i|α⟩(β|a_i⟩ = ⟨a_i|α⟩(a_j|β)*. For spin-1/2 in the $|\pm z\rangle$ basis, $\langle +|S_x=\hbar/2\rangle = 1$, $\langle -|S_x=\hbar/2\rangle = 0$, and, using (1.4.17a), $(\pm |S_x = \hbar/2) = 1/\sqrt{2}$. Therefore

$$|S_x = \hbar/2\rangle\langle S_x = \hbar/2| \doteq \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 0 & 0 \end{pmatrix}$$

 A[|i⟩ + |j⟩] = a_i|i⟩ + a_j|j⟩ ≠ [|i⟩ + |j⟩] so in general it is not an eigenvector, unless a_i = a_j. That is, $|i\rangle + |j\rangle$ is not an eigenvector of A unless the eigenvalues are degenerate.

Sakurai Quantum Mechanics 2nd Edition Instructor Manual

Katrin Zwirglmaier

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